

B] High injection effect:

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- The ambipolar transport equation that we have used to determine the minority charge concentration assumed Low injection.
- As V_{BE} rises, the injected minority carrier concentration may approach, or even become larger than, the majority carrier concentration.
- If we assume quasi-charge neutrality (Remember, that $\int \rho \approx 0 \Rightarrow \int \rho_p = \int \rho_n$), then the

majority carrier hole concentration in p-type base at $x=0$ will rise in fig(a), becoz of excess holes.

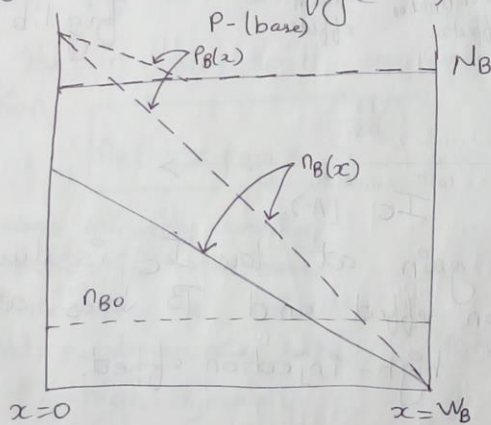


fig 1.a: Minority and majority carrier concentrations in the base under low and high injection.

⇒ Two effects occurs in the BJT at high injection:

- The 1st effect is a reduction in emitter injection efficiency.

• Since, the majority carrier hole concentration at $x=0$ rises with high injection (more holes are injected back into emitter base of F.B. V_{BE}).

• An rise in hole injection causes an rise in I_{EP} current, & an rise in I_{EP} reduces the emitter injection efficiency

$$\beta = \frac{I_{EN}}{I_{EN} + I_{EP}} \quad \beta = \frac{1}{1 + \frac{I_{EP}}{I_{EN}}}$$

b) The common emitter current gain β decreases, with high injection. (β)

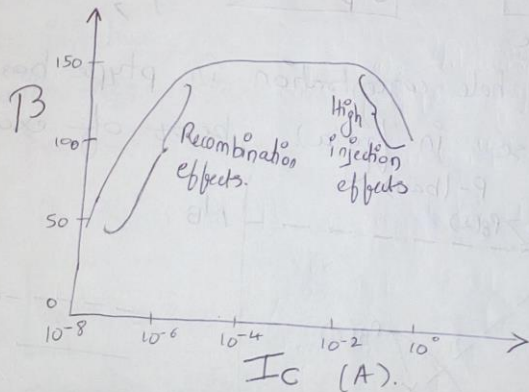


fig 1. b β vs I_C

• The low current gain at low I_C is due to small recombination effect and β decreases at high I_C is due to high-injection effect.

- At low injection, the majority carrier hole concentration is at $x=0$ for npn BJT is,

$$P_B(0) = P_{B0} = N_B$$

- and the minority carrier electron concentration is,

$$n_B(0) = n_{B0} e^{\left(\frac{V_{BE}}{V_T}\right)}$$

- The pn product is,

$$P_B(0) n_B(0) = P_{B0} n_{B0} e^{\left(\frac{V_{BE}}{V_T}\right)} \quad \text{--- (1)}$$

- At high injection, eqⁿ (1) still applies. However, $P_B(0)$ will also rise, and for very high injection it will rise at nearly the same rate as $n_B(0)$.

- The rise in $n_B(0)$ will asymptotically approach the function,

$$n_B(0) \approx n_{B0} e^{\left(\frac{V_{BE}}{2V_T}\right)} \quad \text{--- (2)}$$

- The excess minority carrier concentration in the base and hence the I_C , will rise at a slower rate with V_{BE} in high injection than low injection.

This effect is shown in fig 1.c.

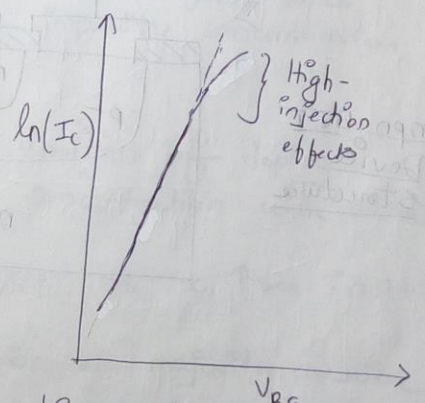


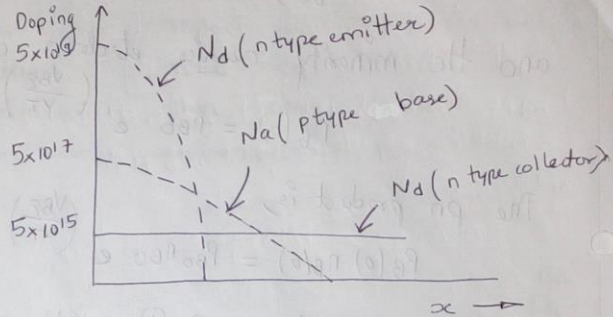
fig 1.c: $\ln(I_C)$ vs V_{BE}

Simulation Expt. (Practice Layouts in Lab): -
 ...

C] Non-uniform Base doping: 20

In analysis of BJT, we assumed uniformly doped regions. However, uniform doping rarely occurs.

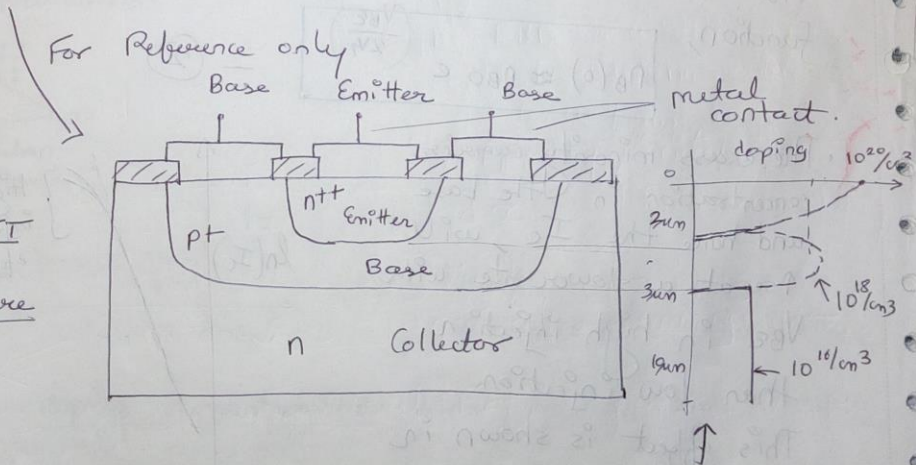
For understanding (Extra):
 (we can start with a uniformly doped ntype substrate, diffuse acceptor atoms from the surface to a compensated ptype base, & then diffuse donor atoms from the surface to form a doubly compensated ntype emitter. The diffusion process results in a non-uniform doping profile).



fig(1.a): Impurity concentration profiles of a double-diffused npn BJT

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n-p-n BJT Device Structure



Doping profile (highly non-uniform)

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Tools: Microwind/Electronic

Extra!

21R

Q. Non-uniform impurity concentration leads to an induced electric field?

⇒ Consider a non-uniformly doped semiconductor with donor impurity atoms. If the semiconductor is in thermal equilibrium, then Fermi-level is constant throughout the semiconductor as shown in fig(b).

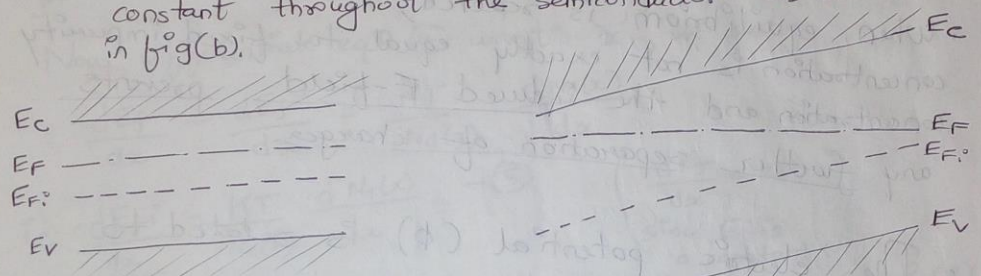


fig a: EBD of a uniformly doped semiconductor.

fig b: EBD for a semiconductor in equilibrium with a nonuniform donor impurity concentration.

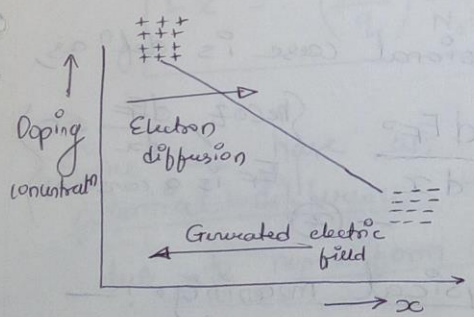


fig c: Variation of doping concentration with distance.

- The doping concentration ↓ as x ↑ in this case.
- ie → There will be a diffusion of majority carriers e^- s from the region of high concentration to

Simulation Expt. (Practice Layouts in Lab): -
Topic: Microscopic/Electro

Extra

region of low concentration, which is in +x direction 22R

- The flow of e^- s leaves behind +vely charged donor ions.
- The separation of +ve and -ve charge induces an E-field i.e. directed to oppose the diffusion process
- When equilibrium is reached, the mobile carrier concentration is not exactly equal to fixed impurity concentration and the induced E-field prevents any further separation of charges.

• The electric potential (ϕ) is related to electron potential energy by charge ($-q$),

$$\phi = \frac{1}{q} (E_F - E_{Fi}) \quad \text{--- (1)}$$

The E-field for one-dimensional case is def^d as,

$$E_x = -\frac{d\phi}{dx} = \frac{1}{q} \frac{dE_{Fi}}{dx} \quad \left\{ \begin{array}{l} \text{because, } \frac{dE_F}{dx} = 0 \\ E_F \text{ is a constant} \end{array} \right. \quad \text{--- (2)}$$

Rule: \leftarrow Equation (2)'s physical meaning: -

* [If intrinsic Fermi level (E_{Fi}) changes as a function of distance through a semiconductor in thermal equilibrium, an Electric field exists in the semiconductor]

Tools: Microminivolt (Practice Layouts in Lab):-

- If we assume a quasi-neutrality condition (in which the electron concentration is almost equal to donor impurity concentration), then

$$n_0 \approx n_i \exp\left(\frac{E_f - E_{f_i}}{kT}\right) \approx N_d(x) \quad \text{--- (3)}$$

Solving for $E_f - E_{f_i} \Rightarrow E_f - E_{f_i} = kT \ln\left[\frac{N_d(x)}{n_i}\right] \quad \text{--- (4)}$

Now differentiate eqⁿ (4) w.r.t 'x', we get

$$\frac{dE_f}{dx} - \frac{dE_{f_i}}{dx} = \frac{kT}{N_d(x)} \frac{dN_d(x)}{dx}$$

i.e. $-\frac{dE_{f_i}}{dx} = \frac{kT}{N_d(x)} \frac{dN_d(x)}{dx} \quad \text{--- (5)}$ $\left(\begin{array}{l} \frac{dE_f}{dx} \approx 0 \\ \text{since } E_f \text{ is a constant} \end{array}\right)$

Thus, E-field can be written, combining eqⁿ (5) & (2),

$$E_x = -\left(\frac{kT}{q}\right) \frac{1}{N_d(x)} \frac{dN_d(x)}{dx} \quad \text{--- (6)}$$

\Rightarrow Since we have an E-field, there will be a potential difference through the semiconductor due to non-uniform doping.

\Downarrow This Result we are gonna use in Non-uniform base doping concept!

Tools: Microwave/Electronic.

Tool: a TCAD Lab on nanohub.org

nanahub.org

FET Lay, CNT Mobility, NanomOS

* Non-uniform base doping ----- (continue) 24

For a p-type base region (npn BJT) in thermal equilibrium, we can write

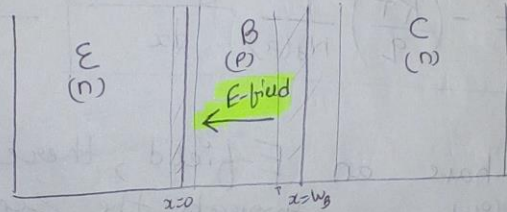
$$J_p = q \mu_p N_a E - q D_p \frac{dN_a}{dx} = 0 \quad \text{--- (i)}$$

Hole current density

$$\text{then, } E = \left(\frac{kT}{q} \right) \frac{1}{N_a} \frac{dN_a}{dx} \quad \text{--- (ii)}$$

(Refer earlier eqn (6))

According to the example of fig (1.a), $\frac{dN_a}{dx}$ is -ve, hence due to non-uniform doping in base, an induced E-field is in -ve x-direction.



Electrons are injected from n-type emitter into the base and the minority carrier base electrons begin diffusing towards the collector region.

The induced E-field in the base, produces force on the electrons in the direction towards the collector.

• The induced E -field, then AIDS the flow of 25 minority carriers across the base region. This E -field is called "accelerating E -field".

• This E -field will produce a drift component of current i.e. in addition to existing diffusion current.

• Since minority carrier electron concentration varies across the base, drift current density will not be a constant.

• The induced electric field in the base due to non-uniform base doping will alter the minority carrier distribution through the base, such that the sum of drift and diffusion current will be a constant.

