

## D] Emitter Bandgap narrowing:

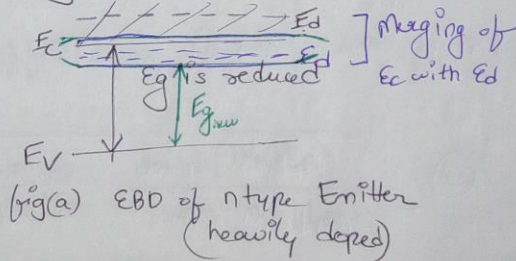
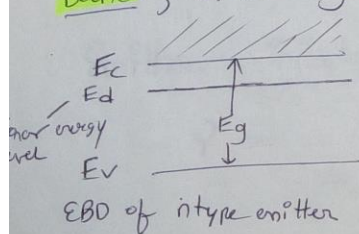
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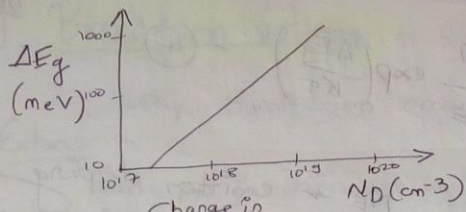
- Another phenomenon affecting the emitter injection efficiency ( $\gamma$ ) is "bandgap narrowing".

$$\gamma \approx \frac{1}{1 + \frac{N_B D_E W_B}{N_E D_B W_E}} \quad \text{--- (1)}$$

From eq<sup>n</sup>, it is clear that  $\gamma$  approaches unity as ratio of base doping to emitter doping continues to decrease as  $\left(\frac{N_B}{N_E}\right)$  decreases.  $\Rightarrow \gamma \rightarrow 1$ .

- As Silicon becomes heavily doped, the discrete donor energy level in an n-type-emitter splits into a band of energies.
- The distance between donor atoms decreases as the concentration of impurity donor atom increases and the splitting of donor level is caused by the interaction of donor atoms with each other.
- As doping  $\downarrow$  (emitter) continues to increase, the donor band widens and moves up toward the conduction band, eventually merging with it. (as shown in fig(a)).





fig(b). Bandgap-energy versus donor impurity concentration in Silicon.

Fig(b) shows a plot of change in bandgap-energy  $\Delta E_g$  versus  $N_D$ .

• A reduction in the  $E_g$  rises the intrinsic carrier concentration.

• The intrinsic carrier concentration is given by,

$$n_i^2 = N_c N_v \exp\left(\frac{-E_{g0}}{KT}\right) \quad \text{--- (1)}$$

$E_{g0} = E_c - E_v$

band-gap  $\rightarrow$  low doping

• In a heavily doped emitter, the  $n_i$  will become.

$$n_{iE}^2 = N_c N_v \exp\left(\frac{-(E_{g0} - \Delta E_g)}{KT}\right)$$

$n_{iE}^2 = n_i^2 \exp\left(\frac{\Delta E_g}{KT}\right)$

--- (2)

where,  $E_{g0}$  is band-gap energy at a low doping and  $\Delta E_g$  is bandgap narrowing factor.

Emitter injector efficiency factor was given by,

$$\gamma = \frac{1}{1 + \frac{p_{E0} D_E L_B}{n_{B0} D_B L_E} \cdot \frac{\tanh(W_B/L_B)}{\tanh(W_E/L_E)}} \quad \text{--- (3)}$$

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$$P_{EO} \approx \frac{n_{iE}^2}{N_E} = \frac{n_i^2}{N_E} \exp\left(\frac{\Delta E_g}{kT}\right) \quad \text{--- (4)}$$

$\downarrow$   
 equilibrium  
 minority carrier  
 concentration in the emitter.

,  $N_E \rightarrow$  emitter doping

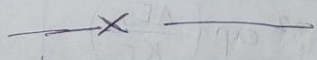
As  $N_E$   $\uparrow$ s,  $\Delta E_g$   $\uparrow$ s, thus it causes  $P_{EO}$  to  $\uparrow$ s (becoz of bandgap narrowing)

As  $P_{EO}$   $\uparrow$ s,  $\gamma$   $\downarrow$ s, this then causes the transistor gain to  $\downarrow$ s.

$\gamma \rightarrow$  emitter injection efficiency.

Thus, a very high emitter doping may result in a smaller current gain that anticipated because of band-gap narrowing effect.

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## E] Breakdown voltage:

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Why Breakdown occurs?

Extra:

From pn junction theory, we know that as you go on rising R.B voltage across a junction, impact ionization within depletion layer sets in and because of impact ionization (there is generation of extra EHPs), reverse current that is flowing through depletion layer starts getting multiplied (EHP generation process). As a result, there is a rapid rise in reverse current in the device, & this is the reason for breakdown.

Qualitatively: Breakdown means  $\Rightarrow$  (sudden/rapid) change in current.

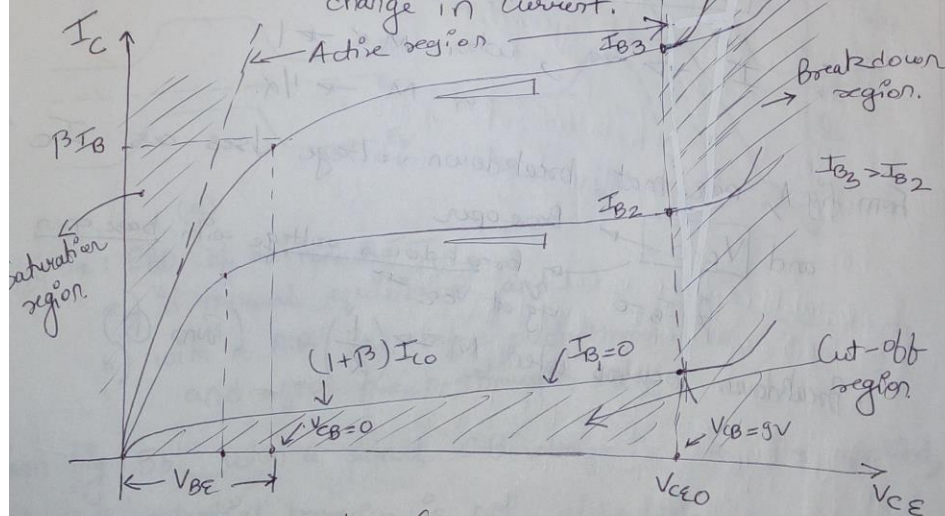


Fig (A)  $I_C$  vs  $V_{CE}$  Curve (npn TR).

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## E] Breakdown voltage:

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• There are two breakdown mechanisms in a BJT

- 1) Punch-through breakdown
- 2) Avalanche breakdown.

### ⇒ Punch-through breakdown:-

As R.B.  $V_{BC}$  ↑, the B-C depletion region widens and extends further into the neutral base.

• It is possible for the B-C depletion region to penetrate completely through the base and reach the B-E depletion region, the effect is called "punch-through".

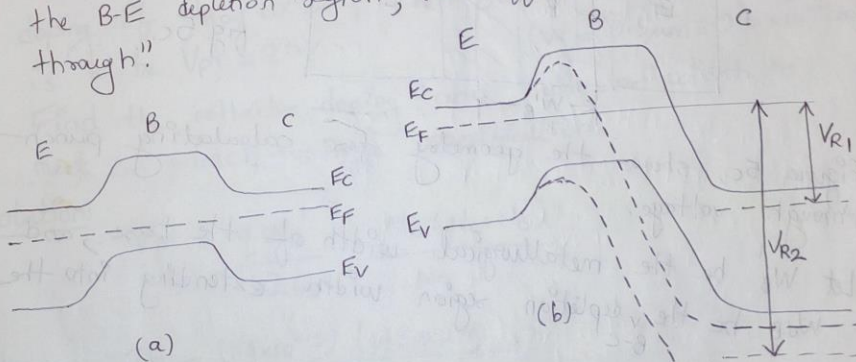


Fig 5a: EBD of an npn BJT

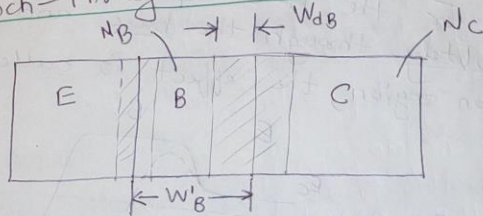
- a) in thermal equilibrium and
- b) with a R.B.  $V_{BC}$  before punch-through  $V_{R1}$  and after punch-through  $V_{R2}$

From fig 5b, when a small C-B voltage i.e.  $V_{R1}$  is applied, the B-E potential barrier is not affected.

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• When a large reverse-bias voltage ( $V_{R2}$ ) is applied, the B-C depletion region extends through the base region and B-E junction potential barrier is lowered because of CB voltage.

• The lowering of potential barrier at BE junction produces a large rise in current, with a very small rise in CB voltage. This effect is called the "punch-through breakdown" phenomenon.



• Figure 5c, shows the geometry for calculating punch-through voltage.

Let  $W_B'$  be the metallurgical width of the base, and  $W_{dB}$  be the depletion region width extending into the base.

(If we neglect narrow depletion width of a F.B, B-E junction, then punch-through occurs when

$$W_{dB} = W_B'$$

$$W_B' = W_{dB} = \left[ \frac{2 \epsilon_s (V_{bi} + V_{pt})}{q} \frac{N_c}{N_B (N_c + N_B)} \right]^{1/2} \quad \text{--- (1)}$$

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where,  $V_{pt}$  is R.B, B-C voltage at punch-through.

If we neglect  $V_{bi} \ll V_{pt}$ , and solve for  $V_{pt}$  from equation (1), we get

$$W_B^2 = \frac{2\epsilon_s V_{pt}}{q} \frac{N_c}{N_B} \left( \frac{1}{N_c + N_B} \right)$$

$$\text{ie } \boxed{V_{pt} = \frac{q W_B^2 N_B (N_c + N_B)}{2\epsilon_s N_c}} \quad \text{--- (2)}$$

Ex. Consider a uniformly doped Si BJT with a base width  $W_B = 0.5 \mu\text{m}$  and a base